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Unconventional gas-phase synthesis of biphenyl and its atropisomeric methyl-substituted derivatives†

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The biphenyl molecule $(C_{12}H_{10})$ acts as a fundamental molecular backbone in the stereoselective synthesis of organic materials due to its inherent twist angle causing atropisomerism in substituted derivatives and in molecular mass growth processes in circumstellar environments and combustion systems. Here, we reveal an unconventional low-temperature phenylethynyl addition–cyclization– aromatization mechanism for the gas-phase preparation of biphenyl $(C_{12}H_{10})$ along with ortho-, meta-, and para-substituted methylbiphenyl $(C_{13}H_{12})$ derivatives through crossed molecular beams and computational studies providing compelling evidence on their formation via bimolecular gas-phase reactions of phenylethynyl radicals (C_6H_5CC, X^2A_1) with 1,3-butadiene- d_6 (C₄D₆), isoprene (CH₂C(CH₃)CHCH₂), and 1,3-pentadiene (CH₂CHCHCHCH₃). The dynamics involve de-facto barrierless phenylethynyl radical additions via submerged barriers followed by facile cyclization and hydrogen shift prior to hydrogen atom emission and aromatization to racemic mixtures (ortho, meta) of biphenyls in overall exoergic reactions. These findings not only challenge our current perception of biphenyls as high temperature markers in combustion systems and astrophysical environments, but also identify biphenyls as fundamental building blocks of complex polycyclic aromatic hydrocarbons (PAHs) such as coronene $(C_{24}H_{12})$ eventually leading to carbonaceous nanoparticles (soot, grains) in combustion systems and in deep space thus affording critical insight into the low-temperature hydrocarbon chemistry in our universe.

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Introduction

Following the discovery of its non-planarity in $1949¹$, the biphenyl $(C_{12}H_{10})$ molecule—a hydrocarbon consisting of two benzene rings bridged by a carbon–carbon single bond—has garnered substantial attention from the organic synthetic, physical organic chemistry, materials science, and medicinal chemistry communities for its role as a vital molecular building block of natural products,² chiroptical nanomaterials,^{3,4} enantioselective catalysts,^{5,6} and photoswitches⁷ (Fig. 1). This significance stems from steric hindrance prohibiting free rotation around the bridging single bond between the phenyl

moieties thus producing a dihedral angle of 44° and rotational barrier of 6 kJ mol⁻¹.^{8,9} The stability of the twisted structure arises from a balance between the steric repulsion of opposing ortho hydrogens and the electrostatic interaction of the π -orbitals.¹⁰ Substituting a single hydrogen of biphenyl induces axial chirality of the molecule and atropisomerism—two key traits for stereoselectivity.^{11,12} Ortho substituents induce the greatest effect, as a replacement of a single hydrogen by a methyl moiety $(-CH_3)$ increases the phenyl-phenyl dihedral angle by 25° , while methyl substitution in the *meta* position offers little to no change in the angle. $13-15$ These properties afford a precise tunability of the biphenyl backbone thus providing unique characteristics and functions which make biaryls vital components of organo- and enantioselective catalysts used in, e.g., asymmetric counteranion-directed catalysis (ACDC) toward the synthesis of chiral drugs and assembly of enantiomerically pure ligands.^{5,16}

The aromatic biphenyl molecule is also of fundamental interest to the astrochemistry and combustion science communities as a precursor to complex polycyclic aromatic

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Fig. 1 Molecular structures of biphenyl [1] and 2-methylbiphenyl [2] as well as examples of a chiral phosphoric acid (CPA) catalyst [3] and a 1,1'binaphthyl-2,2'-diol (BINOL) nanographene [4] emphasizing the twisted biphenyl backbone. Carbon atoms in the biphenyl moiety are colored as black, other carbons are gray, oxygen is red, phosphorus is magenta, hydrogen is white, and side groups are green.

hydrocarbons (PAHs) via the hydrogen abstraction–acetylene addition $(HACA)^{17}$ and phenyl addition-dehydrocyclization $(PAC)^{18}$ mechanisms forming, e.g., phenanthrene $(C_{14}H_{10})$ and triphenylene $(C_{18}H_{12})$, respectively, and eventually carbonaceous nanoparticles referred to as interstellar grains¹⁹ and soot.²⁰ In deep space, PAHs have been suggested to account for up to 30% of the cosmic carbon budget and are implicated as carriers of diffuse interstellar bands $(DIBs)^{21}$ and unidentified infrared (UIR) bands.²² Sophisticated analyses of carbonaceous chondrites such as Allende²³ and Murchison²⁴ revealed the presence of PAHs, including biphenyl, presumably formed in circumstellar envelopes of carbon-rich asymptotic giant branch (AGB) stars along with planetary nebulae as their descendants.²⁵ The formation mechanisms in these high temperature environments are similar to those in hydrocarbon-rich combustion flames, where biphenyl has been observed in flames of alkanes, such as propane,²⁶ *n*-butane,²⁷ and *n*-dodecane,²⁸ alkenes, *e.g.* ethylene,^{29,30} as well as aromatics like benzene,^{31–33} toluene,³⁴ and styrene.³⁵ Traditionally, a bottom-up synthesis of biphenyl proceeds through the reaction of phenyl radicals (C_6H_5) with benzene $(C_6H_6)^{36}$ via an entrance barrier of 11 kJ mol⁻¹.³⁷ Consequently, the aforementioned gas-phase preparation of biphenyl is constrained to high temperature environments where entrance barriers can be overcome easily. While this results in biphenyl playing a vital role in combustion³⁸⁻⁴³ and circumstellar⁴⁴⁻⁴⁶ models for molecular mass growth processes of PAHs at elevated temperatures, pathways to biphenyl in low-temperature environments such as in hydrocarbon rich atmospheres of planets and their moons like Titan and in cold molecular clouds such as TMC-1 and OMC-1 have been discounted and hence remained elusive to date. These pathways are critical to prepare a population of aromatics in low-temperature environments to quantitatively replicate the observed PAH abundances in deep space.⁴⁷⁻⁴⁹

Herein, we report on the gas-phase formation of biphenyl $(C_{12}H_{10})$ and its ortho-, meta-, and para-substituted methylbiphenyl $(C_{13}H_{12})$ derivatives via the bimolecular reactions of the phenylethynyl radical (C $_{6}$ H $_{5}$ CC, X 2 A $_{1})$ with 1,3-butadiene- d_{6} (C_4D_6) , isoprene (2-methyl-1,3-butadiene, $CH_2C(CH_3)CHCH_2$), and $1,3$ -pentadiene (1-methyl-1,3-butadiene, $CH_2CHCHCH$ -CH3) under single collision conditions utilizing a crossed molecular beams machine coupled with high level ab initio electronic structure and statistical calculations (reactions (1)–(3)).

Each elementary reaction proceeds de-facto without an entrance barrier through a van-der-Waals complex in the entrance channel with all transition states leading eventually to biphenyls residing lower in energy than the separated reactants; hence, these systems are rapidly accessible in low-temperature environments such as cold molecular clouds where temperatures as low as 10 K reside.⁵⁰ This pathway constitutes a unique, low-temperature framework for the formation of (substituted) biphenyls and PAHs as racemic building blocks in molecular mass growth processes to carbonaceous nanostructures, while challenging conventional wisdom of biphenyls as high temperature markers in combustion and circumstellar chemistry.

$$
C_6H_5CC (101 amu) + C_4D_6 (60 amu)
$$

\n
$$
\rightarrow C_{12}H_4D_6 (160 amu) + H (1 amu)
$$
 (1a)
\n
$$
\rightarrow C_{12}H_5D_5 (159 amu) + D (2 amu)
$$
 (1b)
\n
$$
C_6H_5CC (101 amu) + CH_2C(CH_3)CHCH_2 (68 amu)
$$

\n
$$
\rightarrow C_{13}H_{12} (168 amu) + H (1 amu)
$$
 (2)

$$
C_6H_5CC
$$
 (101 amu) + CH₂CHCHCHCHCH₃ (68 amu)

$$
\rightarrow C_{13}H_{12}(168 \text{ amu}) + H(1 \text{ amu}) \tag{3}
$$

Methods

Experimental methods

Reactions of phenylethynyl radicals (C_6H_5CC, X^2A_1) with 1,3butadiene- d_6 (C₄D₆, 98.9% D atom, CDN Isotopes), isoprene (CH2C(CH3)CHCH2, 99%, Sigma-Aldrich), and 1,3-pentadiene (CH2CHCHCHCH3, 96%, TCI Chemicals) were conducted under single collision conditions using the crossed molecular beams technique. 51 Each reaction was done at two different collision energies by using either helium (He, 99.9999%, Matheson) or neon (Ne, 99.9999%, Matheson) as a seeding gas for the phenylethynyl beam. Briefly, a (2-bromoethynyl)benzene (C_6H_5CCBr) precursor was purified by multiple freeze–pump–thaw cycles and placed inside the primary source chamber where it was seeded at a fraction of 0.5% in helium or neon at a backing pressure of 500 Torr. The ensuing gas mixture was fed through a Proch-Trickl pulsed valve⁵² operating at -450 V and 120 Hz giving

open times of 80 ms. The precursor was then photodissociated by 193 nm photons from an excimer laser (Coherent, COMPex 110) at 60 Hz and 10 mJ per pulse, creating a phenylethynyl radical beam which passed through a skimmer and was velocity selected by a chopper wheel resulting in a peak velocity (v_p) of 1765 \pm 26 and 896 \pm 10 m s $^{-1}$ and a speed ratio (*S*) of 9.8 \pm 1.1 and 14.0 \pm 1.7 for the helium-seeded and neon-seeded reactions, respectively. The phenylethynyl beam crossed perpendicularly with pulsed beams of 1,3-butadiene- d_6 (v_p = 750 \pm 10 m s $^{-1}$, *S* = 8.0 \pm 0.2), isoprene $(\mathcal{V}_{\rm p}$ = 721 \pm 20 m s $^{-1}$, *S* = 8.5 \pm 0.6), and 1,3-pentadiene ($\mathcal{V}_{\rm p}$ = 711 \pm 20 m s⁻¹, S = 8.5 \pm 0.7) at backing pressures of 550, 450, and 350 Torr, respectively, and pulsed valve parameters of -400 V, 120 Hz, and 80 us open times. These reactions resulted in collision energies ($E_{\rm C}$) of 68.9 \pm 1.5, 72.7 \pm 1.3, and 74.5 \pm 1.5 kJ mol $^{-1}$ and center-of-mass (CM) angles (Θ_{CM}) of 15.0 \pm 0.4, 16.3 \pm 0.5, and $15.8 \pm 0.5^{\circ}$, respectively, for the helium-seeded reactions, while the neon-seeded cases gave E_C values of 25.8 \pm 0.5, 26.8 \pm 0.6, and 26.7 \pm 0.7 kJ mol $^{-1}$ and $\varTheta_{\rm CM}$ values of 27.1 \pm 0.5, 29.2 \pm 0.6, and $28.8 \pm 0.9^{\circ}$ (Table S1, ESI†).

Reactive scattering products were collected by a triply differentially pumped mass spectrometric detector which is rotatable within the plane defined by both reactant beams. Neutral products were ionized at 80 eV and an emission current of 2 mA with a Brink-type⁵³ electron impact ionizer. The consequential ions were filtered by a quadrupole mass spectrometer (QMS, Extrel 150QC) in the time-of-flight (TOF) mode at a constant mass-to-charge ratio (m/z) . Up to 1.8 \times 10⁶ TOF spectra were taken at each angle in 2.5° steps between 9 $\leq \theta \leq 44$ ° with respect to the phenylethynyl beam ($\Theta = 0^{\circ}$). Laboratory angular distributions for each system were created by integrating the TOF spectra and normalizing to the CM angles. To gain information on the reaction dynamics, a forward convolution routine creating user-defined CM translational energy $(P(E_T))$ and angular $(T(\theta))$ flux distributions was used to fit the laboratory angular distributions and TOF spectra in an iterative process.54,55 These functions were used to develop a flux contour map, shown as $I(u, \theta) \sim P(u) \times T(\theta)$,⁵⁶ which portrays an overall image of the outcome of the reaction.

Computational methods

The H loss channels of the reactions of phenylethynyl radicals (C_6H_5CC, X^2A_1) with 1,3-butadiene- d_6 (C_4D_6) , isoprene $(CH_2C(CH_3)CHCH_2)$, and 1,3-pentadiene $(CH_2CHCHCHCH_3)$ on adiabatic doublet ground-state potential energy surfaces were explored, along with the phenylacetylene (C_6H_5CCD) formation path in the phenylethynyl-1,3-butadiene- d_6 system. The geometries of collision complexes, cyclized and hydrogenshifted intermediates, transition states, and H loss products, were optimized with density functional w B97X-D 57 /cc-pVTZ calculations. Their coupled cluster $CCSD(T)^{58-60}/cc$ -pVTZ energies with wB97X-D/cc-pVTZ zero-point energy correction were then obtained. Single-point higher level (multi-configuration) energy has been shown to yield good results for various closed and open-shell systems, $61-65$ where the methods of CCSD(T)// B3LYP, G3//B3LYP, G3//B3LYP, and G3//wB97X-D, respectively, were employed. Intrinsic reaction coordinate (IRC) calculations have been carried out for the collision complexes **i1**, **i5**, **i9**, **i13**, and i17—the structures of which are discussed in the Potential energy surfaces section—back to reactants by density functional wB97X-D. GAUSSIAN16⁶⁶ programs were facilitated in the *ab initio* electronic structure calculations. The energy dependent RRKM (Rice–Ramsperger–Kassel–Marcus) rate constants⁶⁷ were computed for the likely atomic hydrogen loss products—including 4- 5- and 6-membered ring species—on doublet surfaces at various experimental collision energies and at 0.0 kJ mol^{-1} . The species involved were treated as a collection of harmonic oscillators. The saddle-point approach^{67,68} was utilized to estimate the number of states of transition states, and density of states of collision complexes and intermediates with wB97X-D/cc-pVTZ harmonic frequencies and CCSD(T)/cc-pVTZ energies. In order to obtain branching ratios for all accessible products, complete reaction paths are required. Preliminary calculations illustrate, for instance, there are at least 16 paths for the collision complex i5 (detailed in the Potential energy surfaces section) to isomerize, which means it is not practical to locate complete paths. However, it is feasible to estimate branching ratios among products by comparing the bottleneck (smallest) rate constant along p1, p2, and p10 formation paths, which strongly indicates p1 (biphenyl) is the major product. In order to get a quantitative result, the product branching ratios at zero collision energy were computed by solving the rate equations derived from *ab initio* reaction paths as in Fig. S4–S6 (ESI†), respectively, with the Runge–Kutta method. The concentration evolution with time was thus obtained for every species involved, in which the asymptotic value of concentration is taken as branching ratio with collision complex initial concentration assumed to be 1.

Results & discussion

Laboratory frame

Reactive scattering time-of-flight (TOF) spectra (Fig. 2) were obtained at two different collision energies for the reactions of phenylethynyl (C₆H₅CC) radicals with 1,3-butadiene- d_6 (C₄D₆), isoprene $(CH_2C(CH_3)CHCH_2)$, and 1,3-pentadiene $(CH_2CH CHCHCH₃$). First, signal was observed for the bimolecular reaction of phenylethynyl (C_6H_5CC) with 1,3-butadiene- d_6 (C_4D_6) at mass-to-charge ratios (m/z) of 160 $(C_{12}H_4D_6^+, {^{13}CC_{11}}$ $H_5D_5^+$, ${}^{13}CC_{11}H_3D_6^+$), 159 $(C_{12}H_5D_5^+$, $C_{12}H_3D_6^+$, ${}^{13}CC_{11}H_4D_5^+$, ¹³CC₁₁H₂D₆⁺), and 158 (C₁₂H₄D₅⁺, C₁₂H₂D₆⁺, ¹³CC₁₁H₅D₄⁺, ${}^{13}CC_{11}H_{3}D_{5}^{+}$, ${}^{13}CC_{11}HD_{6}^{+}$) for both high $(E_{C}$ = 68.9 \pm 1.5 kJ mol⁻¹) and low ($E_C = 25.8 \pm 0.5$ kJ mol⁻¹) collision energies. In both cases, the TOFs at distinct m/z ratios overlap after scaling (Fig. S1, ESI†) with $m/z = 160$ collected at a level of 13 \pm 5% compared to m/z = 159; thus, signal at m/z = 160 likely arises from ¹³C-substituted $C_{12}H_5D_5^+$ accounting for natural isotopic abundance of about 1.1% for ¹³C, while that at $m/z =$ 159 originates from $C_{12}H_5D_5^+$ coupled with atomic deuterium loss (reaction (1b)). This verifies that the hydrogen (deuterium) loss occurs solely from the 1,3-butadiene reactant. Since electron impact ionization at 80 eV was applied to ionize reactively

Fig. 2 Laboratory angular distributions (LADs) (a)–(c), (g)–(i) and time-of-flight spectra (TOFs) (d)–(f), (j)–(l) recorded at mass-to-charge (m/z) = 159, 168, and 168 for the reactions of phenylethynyl (C₆H₅CC) with 1,3-butadiene- d_6 (C₄D₆, (a), (d), (g), (j)), isoprene (CH₂C(CH₃)CHCH₂, (b), (e), (h), (k)), and 1,3pentadiene (CH₂CHCHCHCH₃, (c), (f), (i), (l)), respectively. CM represents the center-of-mass angle, and 0° and 90° define the directions of the phenylethynyl and 1,3-butadiene- d_6 /isoprene/1,3-pentadiene beams, respectively. The black circles depict the data and red lines the fits. Reactions were done with high (a)–(f) and low (g)–(l) collision energies. Carbon atoms are colored gray, hydrogens are white, and deuterium atoms are light blue.

scattered products, the counts observed at $m/z = 158$ result from dissociative electron impact ionization of the parent $C_{12}H_5D_5$ species in the ionizer. TOF spectra were then collected at $m/z =$ 159 at center-of-mass (CM) angles (Θ_{CM}) of 14.3 \pm 0.4 (Fig. 2d) and 26.3 \pm 0.5° (Fig. 2j) with respect to the phenylethynyl beam for high and low collision energies, respectively. The high E_{C} TOFs are very narrow, spanning only about 150 µs, while the low E_C TOFs are much broader, spanning about 300 μ s. TOFs were also acquired at higher and lower angles in 2.5° steps in

order to create laboratory angular distributions (LAD, Fig. 2a and g), which are nearly symmetric about Θ_{CM} , indicating indirect reaction dynamics through $C_{12}H_5D_6$ intermediate(s). A strongly forward or backward peaking of the LAD with respect to Θ_{CM} would indicate rather direct scattering dynamics, which were not observed here.

For both high (72.7 \pm 1.3 kJ mol⁻¹) and low (26.8 \pm 0.6 kJ mol $^{-1}$) collision energies in the reaction of phenylethynyl (C_6H_5CC) with isoprene $(CH_2C(CH_3)CHCH_2)$, signal was

observed at $m/z = 169 \; (^{13} \text{CC}_{12} \text{H}_{12}^{\dagger})$, 168 $(\text{C}_{13} \text{H}_{12}^{\dagger}, \; ^{13} \text{CC}_{12} \text{H}_{11}^{\dagger})$, and 167 $(C_{13}H_{11}^{\text{+}},{}^{13}CC_{12}H_{10}^{\text{+}})$. These peaks overlap after scaling revealing that all three masses arise from a single channel (Fig. S2, ESI†). Signal at $m/z = 169$ shows a level of 12 \pm 5% compared to $m/z = 168$ which is indicative of natural ¹³C-substituted $C_{13}H_{12}$ ⁺ species, while the peak at $m/z = 168$

Fig. 3 CM product translational energy (a), (d), (g), (j), (m), (p) and angular (b), (e), (h), (k), (n), (q) flux distributions, as well as the associated flux contour maps (c), (f), (i), (l), (o), (r) leading to the formation of C₁₂H₅D₅, C₁₂H₁₃, and C₁₂H₁₃ isomers in the reactions of phenylethynyl (C₆H₅CC) with 1,3-butadiene d_6 (C₄D₆, (a)–(f)), isoprene (CH₂C(CH₃)CHCH₂, (g)–(l)), and 1,3-pentadiene (CH₂CHCHCHCH₃, (m)–(r)). Red lines define the best-fit functions while shaded areas provide the error limits. The flux contour map represents the intensity of the reactively scattered products as a function of product velocity (u) and scattering angle (θ) , and the color bar indicates flux gradient from high (H) to low (L) intensity. Carbon atoms are colored gray, hydrogens are white, and deuterium atoms are light blue.

originates from ionized $C_{13}H_{12}$ along with atomic hydrogen (reaction (2)). The counts observed at $m/z = 167$ result from dissociative electron impact ionization of the parent $C_{13}H_{12}$ species in the ionizer. Signal for the reaction of phenylethynyl with the 1,3-pentadiene (CH₂CHCHCHCH₃) isomer at E_C = 74.5 \pm 1.5 and 26.7 \pm 0.9 kJ mol⁻¹ was much lower (Fig. S3, ESI[†]) with signal only detected at $m/z = 168$ and 167 (reaction (3)). Background intensity from the (2-bromoethynyl)benzene precursor at $m/z = 154$ prevented detection of reactive scatting signal for the possible methyl loss channel. For both the isoprene and 1,3-pentadiene systems, TOFs were collected at $m/z = 168$ at CM angles of 15.3 \pm 0.5 (Fig. 2e) and 15.3 \pm 0.5° (Fig. 2f) for high E_C and 28.3 \pm 0.6 (Fig. 2k) and 28.3 \pm 0.9° (Fig. 2l) for low E_C , respectively. Much like the 1,3-butadiene- d_6 system, the TOFs for the C_5H_8 isomer reactions at high collision energies are very narrow at about 150 µs wide, while those at low collision energies are about twice the width. The LADs are also quite similar, with forward–backward symmetry about Θ_{CM} implying indirect reaction dynamics leading to $C_{13}H_{12}$ product(s) along with atomic hydrogen through $C_{13}H_{13}$ intermediate(s) (Fig. 2b,c and h,i).

Center-of-mass frame

With the identification of the $C_{12}H_5D_5$ product(s) coupled with deuterium loss from the phenylethynyl-1,3-butadiene- d_6 system, as well as $C_{13}H_{12}$ product(s) via hydrogen loss from the phenylethynyl–isoprene/1,3-pentadiene reactions, we now attempt an identification of the nature of the intermediates and products along with the overall reaction mechanisms. This is achieved by converting the laboratory data (TOF spectra, LAD) to the CM reference frame, 69 producing the CM product translational energy $(P(E_T))$ and angular $(T(\theta))$ flux distributions (Fig. 3) and merging these data with electronic structure calculations (Fig. 4–6). For all three reactions and collision

energies, the laboratory data could be fit with a single channel corresponding to product masses of 159 amu $(C_{12}H_5D_5)$ plus 2 amu (D) (phenylethynyl-1,3-butadiene- d_6) and 168 amu $(C_{13}H_{12})$ plus 1 amu (H) (phenylethynyl-isoprene/1,3-pentadiene).

In case of reaction (1), the $P(E_T)$ extends up to 423 \pm 50 and 382 \pm 54 kJ mol⁻¹ for high and low collision energies, respectively (Fig. 3a and d). Utilizing energy conservation for those products born without internal excitation, given by $\Delta_r G$ = $E_{\rm C}$ – $E_{\rm max}$, the reaction energy $(\Delta_{\rm r}G)$ can be obtained by subtracting the maximum translational energy (E_{max}) from the collision energy ($E_{\rm C}$) providing similar values of $-354~\pm$ 52 and -356 ± 55 kJ mol⁻¹ at high and low collision energies. The $P(E_{\rm T})$ depicts a maximum at about 45 kJ mol $^{-1}$ at both collision energies suggesting that $C_{12}H_5D_5$ product(s) are formed via a tight exit transition state involving substantial electron density rearrangement from the decomposing complex to the final products.⁷⁰ The angular flux distributions show intensity over the full angular range, reinforcing the implication of indirect scattering dynamics through $C_{12}H_5D_5$ intermediate(s) (Fig. 3b and e). The $T(\theta)$ also display a slight forward asymmetry with an intensity ratio $I(0^{\circ})/I(180^{\circ})$ of about (1.6 ± 0.3) : 1.0, which suggests the existence of at least one channel where complex formation takes place but the lifetime is too short to allow multiple rotations.⁷¹ These findings for both collision energies are reflected in the flux contour maps (Fig. 3c and f).

For reaction (2), the $P(E_T)$ provides E_{max} values of 441 \pm 41 and 382 \pm 49 kJ mol⁻¹ yielding reaction energies of -368 ± 42 and -355 ± 50 kJ mol $^{-1}$ for high and low $E_{\rm C}$, respectively, while the $P(E_{\rm T})$ for reaction (3) affords reaction energies of $-380\,\pm\,$ 60 and -358 ± 61 kJ mol⁻¹ (Fig. 3g, j, m, and p). At high collision energy, the $P(E_T)$ for both reactions (2) and (3) peak at about 95 kJ mol $^{-1}$ while the $P(E_{\rm T})$ peaks at about 35 kJ mol $^{-1}$

Fig. 4 Schematic potential energy surface for the reaction of phenylethynyl (C₆H₅CC) with 1,3-butadiene-d₆ (C₄D₆) obtained with CCSD(T)/cc-pVTZ// wB97XD/cc-pVTZ calculations.

at low collision energy; therefore, both reactions (2) and (3) involve a tight exit transition state leading to $C_{13}H_{12}$ product(s) coupled with atomic hydrogen loss. Inspecting the angular flux distributions (Fig. 3h, k, n, and q), the $T(\theta)$ for reactions (2) and (3) feature isotropic scattering in which there is equal intensity at all angles. This reveals the formation of long-lived $C_{13}H_{13}$ intermediate(s) with lifetime(s) longer than their rotational period(s). These results are mirrored in the flux contour maps (Fig. 3i, l, o, and r). Overall, the reactions to form $C_{13}H_{12}$ are strongly exoergic by 355 \pm 50 to 380 \pm 60 kJ mol $^{-1}$.

Potential energy surfaces

The nature of the products and reaction pathways are now elucidated by combining our experimental results with electronic

structure and statistical calculations. These potential energy surfaces (PES) were computed for the intermediates, reactants, and transition states of each system (Fig. 4–6 and Fig. S4–S6, ESI†). Rate constants were calculated exploiting Rice–Ramsperger–Kassel–Marcus (RRKM) theory⁷² at the experimental collision energies, as well as at 0 kJ mol^{-1} to account for interstellar conditions (Tables S2–S4, ESI†). The differences in energetics between deuterated and nondeuterated species are minor.

For reaction (1b), the experiment provides a reaction energy $(-354 \pm 52; -356 \pm 55 \text{ kJ mol}^{-1})$ which coincides with the calculated energy of biphenyl- d_5 (C₁₂H₅D₅, p1) plus atomic deuterium $(-361 \text{ kJ mol}^{-1}, \text{Fig. 4}).$ The initial step of reaction starts with the formation of a van-der-Waals complex (vdw1) in the entrance channel. The latter isomerizes via phenylethynyl

Fig. 5 Schematic potential energy surface for the reaction of phenylethynyl (C₆H₅CC) with isoprene (CH₂C(CH₃)CHCH₂) obtained with CCSD(T)/ccpVTZ//wB97XD/cc-pVTZ calculations for phenylethynyl addition to the C4 (a) and C1 (b) carbons of isoprene.

radical addition to one of the double bonds of 1,3-butadiene- d_6 leading to methylallyl radical collision complex i1, which is stabilized by 282 kJ mol^{-1} with respect to the separated reactants. The barrier to addition of only 1 kJ mol^{-1} is well below the energy of the separated reactants (submerged barrier).

While there are two possible pathways leading to p1, the energetically favorable route involves six-membered ring closure of the radical side chain (i1 \rightarrow i3), [1,2] deuterium shift on the newly formed cyclic group (i3 \rightarrow i4), and finally deuterium loss (i4 \rightarrow p1) over a tight exit barrier. While the experimental reaction energy matches that of p1, it is possible that 3,5 hexadien-1-yn-1-ylbenzene (p2) is veiled within the lower energy

portion of the $P(E_T)$ and acts as a secondary product. There are also two pathways leading to p2, with the most likely being a simple addition–elimination mechanism featuring deuterium loss from methylallyl-type radical i1. The rate constants reinforce these conclusions, where the six-membered ring closure *via* the methylallyl-type radical $i1 \rightarrow$ phenylcyclohexadienyl radical i3 pathway (k_1) , subsequent $[1,2]$ deuterium shift, and atomic deuterium loss is at least three orders of magnitude faster than the alternative reaction routes (k_2, k_3) from the methylallyl-type radical i1 as discussed above, while routes involving the formation of a 5-membered ring are also slower (Table S2 and Fig. S4, ESI†). Overall, these results provide compelling evidence on the formation of at least biphenyl- d_5

Fig. 6 Schematic potential energy surface for the reaction of phenylethynyl (C₆H₅CC) with 1,3-pentadiene (CH₂CHCHCHCH $_3$) obtained with CCSD(T)/ cc-pVTZ//wB97XD/cc-pVTZ calculations for phenylethynyl addition to the C1 (a) and C4 (b) carbons of 1,3-pentadiene.

via the reaction of the phenylethynyl radical with 1,3-butadiene- d_6 .

Reactions (2) and (3) follow a similar pattern as for reaction (1b), with the methyl group on the isoprene and 1,3-pentadiene reactants acting as a spectator, i.e. the methyl moiety is not actively participating in the chemical dynamics of the reactions. By changing the position of the methyl group on the 1,3 butadiene backbone from the C1 to the C2 carbon atom, the resulting methyl position on the biphenyl moiety can be tuned. For reaction (2) (Fig. 5), the phenylethynyl radical can add to either of the terminal double-bonded carbons on the isoprene backbone. The initially formed van-der-Waals complexes (vdW2, vdW3) may isomerize. The phenylethynyl radical addition to the C4 (Fig. 5a) or C1 (Fig. 5b) carbon of isoprene leads eventually to 3-methylbiphenyl $(C_{13}H_{12}$, p3) or 4-methylbiphenyl $(C_{13}H_{12}$, p5), respectively, plus atomic hydrogen through the most favorable pathways $i5 \rightarrow i7 \rightarrow i8 \rightarrow p3$ or $i9 \rightarrow i11 \rightarrow$ $i12 \rightarrow p5$, which are the methyl-substituted equivalents of the phenylethynyl–1,3-butadiene- d_6 route. The calculated reaction energies for ${\bf p}3$ $(-367\ {\rm kJ\ mol}^{-1})$ and ${\bf p5}$ $(-367\ {\rm kJ\ mol}^{-1})$ match nicely with the observed experimental reaction energies $(-368 \pm 42;\ -355 \pm 50 \, \, \mathrm{kJ} \, \, \mathrm{mol}^{-1})$. Likewise, for reaction (3), van-der-Waals complexes (vdW4, vdW5) can be accessed in the entrance channel, and their fate is dictated by an addition of the phenylethynyl radical via a submerged barrier to the doubly-bonded C1 (Fig. 6a) or C4 (Fig. 6b) carbon of

Fig. 7 Product channels for the reactions of phenylethynyl radicals (C₆H₅CC) with 1,3-butadiene- d_6 (C₄D₆), isoprene (CH₂C(CH₃)CHCH₂), and 1,3-pentadiene (CH₂CHCHCHCH₃). Carbon atoms are colored gray, hydrogens are white, and deuteriums are light blue.

1,3-pentadiene leading eventually to 2-methylbiphenyl $(C_{13}H_{12}$, p7) or 3-methylbiphenyl $(C_{13}H_{12}$, p3), respectively, via hydrogen loss through the phenylethynyl–isoprene equivalent $i13 \rightarrow i15 \rightarrow i16 \rightarrow p7$ and $i17 \rightarrow i19 \rightarrow i20 \rightarrow p3$ routes. The rate constants from the initial methylallyl-type radical collision complexes (i5/i9/i13/i17) also support this reasoning, with the

Fig. 8 Key role of methyl substituted biphenyls exploiting 2-methylbiphenyl as a prototype through the hydrogen abstraction–acetylene addition (HACA, red), hydrogen abstraction–vinylacetylene addition (HAVA, blue), and radical–radical reaction (RRR, green) mechanisms to the products 7-phenylindene, 4-methylphenanthrene, 1-phenylnaphthalene, and [4]-helicene, as well as the cis- and trans-1,1'-binaphthyl backbones of compounds [3] and [4] from Fig. 1 eventually leading to coronene—a molecular building block of two-dimensional nanostructures. Carbon atoms are colored gray and hydrogens are white.

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cyclization step followed by [1,2] hydrogen shift and hydrogen atom loss giving rate constant values of up to eight orders of magnitude higher than alternate pathways (Tables S3–S4, ESI†). Additionally, in all three reactions, the critical transition state in the formation of the monocyclic products is $53-73$ kJ mol⁻¹ higher in energy than for the biphenyl derivatives, indicating that the products p2, p4, p6, p8, and p9 encompass only a minor fraction of reactive scattering signal, if any, which further reinforces the generation of biphenyl and its methylsubstituted derivatives as major products in the reactions of phenylethynyl radicals with 1,3-butadiene- d_6 , isoprene, and 1,3-pentadiene (Fig. 7).

Conclusion & outlook

Overall, our combined crossed molecular beams and computational study on the bimolecular gas-phase reactions of phenylethynyl radicals (C_6H_5CC) with three molecules containing a 1,3-butadiene backbone—1,3-butadiene- d_6 (C₄D₆), isoprene $(CH_2C(CH_3)CHCH_2)$, and 1,3-pentadiene $(CH_2CHCHCHCH_3)$ uncovered the unconventional formation of biphenyl and its ortho-, meta-, and para-substituted methylbiphenyl derivatives under single collision conditions. The reaction routes involve formation of van-der-Waals complexes in the entrance channels of the reactions followed by isomerization through addition of the phenylethynyl radical center via submerged barriers to the π -electron system of the secondary reactant at the terminal carbon atom(s) followed by facile ring closure and hydrogen shift isomerization prior to unimolecular decomposition of the reaction intermediates via atomic hydrogen loss. These patterns classify this reaction sequence as a lowtemperature framework for the formation of (un)substituted biphenyls and PAHs as building blocks in molecular mass growth processes to carbonaceous nanostructures. Prototypes of successive reactions, of which the fundamental mechanisms have previously been explored, $73-78$ are featured in Fig. 8 eventually leading to the formation of coronene $(C_{24}H_{12})$ —a fundamental molecular building block in two-dimensional carbonaceous nanostructures—via hydrogen abstraction–acetylene addition (HACA), hydrogen abstraction–vinylacetylene addition (HAVA), and radical–radical reactions (RRR).

Essentially, this phenylethynyl addition–cyclization–aromatization mechanism serves as a versatile framework for tuning the position of an alkyl group within the biphenyl backbone to adjust chirality and dihedral angle for stereoselectivity. Further, the de-facto barrierless nature of reaction defies conventional wisdom that biphenyl symbolizes a high temperature marker in combustion and circumstellar chemistry as a product of, e.g., phenyl radicals with benzene.³⁶ This opens up (un)substituted biphenyl formation through previously overlooked reaction pathways in low-temperature environments such as cold molecular clouds and hydrocarbon-rich atmospheres of planets or their moons. While PAHs like biphenyl itself are difficult to explicitly detect in extraterrestrial environments due to the lack of a permanent dipole moment, ortho-, meta-, and/or

Data availability

The data that support the findings of this study are available in the article and the ESI.† Additional data are available from the corresponding authors upon reasonable request.

Author contributions

R. I. K. designed and supervised the experiment; S. J. G., C. H., Z. Y., and A. S. performed the experimental measurements; A. P.-G. synthesized the experimental precursor; W. S. supervised the precursor synthesis; S. J. G. performed the data analysis; B.-J. S., S. F., and K. P. K. carried out the threoretical analysis; A. H. H. C. supervised the theoretical analysis; S. J. G., R. I. K., W. S., and A. H. H. C. wrote the paper.

Conflicts of interest

The authors declare no competing financial interest.

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References

- 1 O. Bastiansen, Acta Chem. Scand., 1949, 3, 408–414.
- 2 T. Qin, S. L. Skraba-Joiner, Z. G. Khalil, R. P. Johnson, R. J. Capon and J. A. Porco, Nat. Chem., 2015, 7, 234–240.
- 3 S. Li, R. Li, Y.-K. Zhang, S. Wang, B. Ma, B. Zhang and P. An, Chem. Sci., 2023, 14, 3286–3292.
- 4 T. M. Fukunaga, T. Kato, K. Ikemoto and H. Isobe, Proc. Natl. Acad. Sci. U. S. A., 2022, 119, e2120160119.
- 5 L. C. Mayer, S. Heitsch and O. Trapp, Acc. Chem. Res., 2022, 55, 3345–3361.
- 6 T. A. Schmidt, S. Schumann, A. Ostertag and C. Sparr, Angew. Chem., Int. Ed., 2023, e202302084, DOI: [10.1002/](https://doi.org/10.1002/anie.202302084) [anie.202302084](https://doi.org/10.1002/anie.202302084).
- 7 A. H. G. David, R. Casares, J. M. Cuerva, A. G. Campaña and V. Blanco, J. Am. Chem. Soc., 2019, 141, 18064–18074.
- 8 A. Almenningen, O. Bastiansen, L. Fernholt, B. N. Cyvin, S. J. Cyvin and S. Samdal, J. Mol. Struct., 1985, 128, 59–76.
- 9 O. Bastiansen and S. Samdal, J. Mol. Struct., 1985, 128, 115–125.
- 10 K. C. Park, L. R. Dodd, K. Levon and T. K. Kwei, Macromolecules, 1996, 29, 7149–7154.
- 11 G. Bringmann, A. J. Price Mortimer, P. A. Keller, M. J. Gresser, J. Garner and M. Breuning, Angew. Chem., Int. Ed., 2005, 44, 5384–5427.
- 12 C. B. Roos, C.-H. Chiang, L. A. M. Murray, D. Yang, L. Schulert and A. R. H. Narayan, Chem. Rev., 2023, 123, 10641–10727.
- 13 J. P. Maier and D. W. Turner, Faraday Discuss. Chem. Soc., 1972, 54, 149–167.
- 14 J. J. Dynes, F. L. Baudais and R. K. Boyd, Can. J. Chem., 1985, 63, 1292–1299.
- 15 L. Andrews, R. T. Arlinghaus and C. K. Payne, J. Chem. Soc., Faraday Trans. 2, 1983, 79, 885–895.
- 16 M. Mahlau and B. List, Angew. Chem., Int. Ed., 2013, 52, 518–533.
- 17 T. Yang, T. P. Troy, B. Xu, O. Kostko, M. Ahmed, A. M. Mebel and R. I. Kaiser, Angew. Chem., Int. Ed., 2016, 55, 14983–14987.
- 18 L. Zhao, M. B. Prendergast, R. I. Kaiser, B. Xu, U. Ablikim, M. Ahmed, B. J. Sun, Y. L. Chen, A. H. H. Chang and R. K. Mohamed, Angew. Chem., Int. Ed., 2019, 131, 17603–17611.
- 19 B. T. Draine, Annu. Rev. Astron. Astrophys., 2003, 41, 241–289.
- 20 J. W. Martin, M. Salamanca and M. Kraft, Prog. Energy Combust. Sci., 2022, 88, 100956.
- 21 W. W. Duley, Faraday Discuss., 2006, 133, 415–425.
- 22 A. M. Ricks, G. E. Douberly and M. A. Duncan, Astrophys. J., 2009, 702, 301.
- 23 L. Becker and T. E. Bunch, Meteorit. Planet. Sci., 1997, 32, 479–487.
- 24 T. N. Tingle, C. H. Becker and R. Malhotra, Meteoritics, 1991, 26, 117–127.
- 25 F. L. Plows, J. E. Elsila, R. N. Zare and P. R. Buseck, Geochim. Cosmochim. Acta, 2003, 67, 1429–1436.
- 26 N. M. Marinov, M. J. Castaldi, C. F. Melius and W. Tsang, Combust. Sci. Technol., 1997, 128, 295–342.
- 27 N. M. Marinov, W. J. Pitz, C. K. Westbrook, A. M. Vincitore, M. J. Castaldi, S. M. Senkan and C. F. Melius, Combust. Flame, 1998, 114, 192–213.
- 28 T. Mitra, T. Zhang, A. D. Sediako and M. J. Thomson, Combust. Flame, 2019, 202, 33–42.
- 29 R. P. Lindstedt and B. B. O. Waldheim, Proc. Combust. Inst., 2013, 34, 1861–1868.
- 30 A. Liu, C. E. Garcia, F. Sewerin, B. A. O. Williams and S. Rigopoulos, Combust. Flame, 2020, 221, 384–400.
- 31 H. Richter, T. G. Benish, O. A. Mazyar, W. H. Green and J. B. Howard, Proc. Combust. Inst., 2000, 28, 2609–2618.
- 32 L. Zhao, W. Sun, J. Yang and B. Yang, Combust. Flame, 2018, 197, 355–368.
- 33 A. Hamadi, W. Sun, S. Abid, N. Chaumeix and A. Comandini, Combust. Flame, 2022, 237, 111858.
- 34 V. Detilleux and J. Vandooren, J. Phys. Chem. A, 2009, 113, 10913–10922.
- 35 W. Yuan, Y. Li, P. Dagaut, J. Yang and F. Qi, Combust. Flame, 2015, 162, 1868–1883.
- 36 F. Zhang, X. Gu and R. I. Kaiser, J. Chem. Phys., 2008, 128, 084315.
- 37 J. Park, S. Burova, A. S. Rodgers and M. C. Lin, J. Phys. Chem. A, 1999, 103, 9036–9041.
- 38 J. Guo, P. Liu, E. Quadarella, K. Yalamanchi, I. Alsheikh, C. Chu, F. Liu, S. M. Sarathy, W. L. Roberts and H. G. Im, Combust. Flame, 2022, 246, 112420.
- 39 A. Liu, Z. Gao, S. Rigopoulos, K. H. Luo and L. Zhu, Fuel, 2022, 317, 122897.
- 40 T. Kathrotia, P. Oßwald, C. Naumann, S. Richter and M. Köhler, Fuel, 2021, 302, 120736.
- 41 X. Shi, Q. Wang and A. Violi, Combust. Flame, 2020, 212, 216–233.
- 42 T. Zhang and M. J. Thomson, Combust. Flame, 2018, 190, 416–431.
- 43 J.-L. Consalvi, F. Liu, M. Kashif and G. Legros, Combust. Flame, 2017, 180, 167–174.
- 44 D. K. Bohme, Chem. Rev., 1992, 92, 1487–1508.
- 45 I. Cherchneff, in Molecules in Astrophysics: Probes and Processes, ed. E. F. v Dishoeck, Cambridge University Press, 1997, vol. 178, pp. 469–476.
- 46 P. Cau, Astron. Astrophys., 2002, 392, 203–213.
- 47 A. G. G. M. Tielens, Annu. Rev. Astron. Astrophys., 2008, 46, 289–337.
- 48 L. J. Allamandola, D. M. Hudgins and S. A. Sandford, Astrophys. J., 1999, 511, L115.
- 49 L. Zhao, R. I. Kaiser, B. Xu, U. Ablikim, M. Ahmed, M. M. Evseev, E. K. Bashkirov, V. N. Azyazov and A. M. Mebel, Nat. Astron., 2018, 2, 973–979.
- 50 S. T. Megeath, E. Allgaier, E. Young, T. Allen, J. L. Pipher and T. L. Wilson, Astron. J., 2009, 137, 4072.
- 51 X. Gu, Y. Guo and R. I. Kaiser, Int. J. Mass Spectrom., 2005, 246, 29–34.
- 52 D. Proch and T. Trickl, Rev. Sci. Instrum., 1989, 60, 713–716.
- 53 G. O. Brink, Rev. Sci. Instrum., 1966, 37, 857–860.
- 54 M. F. Vernon, PhD Dissertation, University of California at Berkeley, 1983.
- 55 P. S. Weiss, PhD Dissertation, University of California at Berkeley, 1985.
- 56 R. I. Kaiser, Chem. Rev., 2002, 102, 1309–1358.
- 57 J.-D. Chai and M. Head-Gordon, Phys. Chem. Chem. Phys., 2008, 10, 6615–6620.
- 58 G. D. Purvis, III and R. J. Bartlett, J. Chem. Phys., 1982, 76, 1910–1918.
- 59 C. Hampel, K. A. Peterson and H.-J. Werner, Chem. Phys. Lett., 1992, 190, 1–12.
- 60 P. J. Knowles, C. Hampel and H. J. Werner, J. Chem. Phys., 1993, 99, 5219–5227.
- 61 J. M. Simmie and K. P. Somers, J. Phys. Chem. A, 2015, 119, 7235–7246.
- 62 B. M. Jones, F. Zhang, R. I. Kaiser, A. Jamal, A. M. Mebel, M. A. Cordiner and S. B. Charnley, Proc. Natl. Acad. Sci. U. S. A., 2011, 108, 452–457.
- 63 K. L. K. Lee, B. A. McGuire and M. C. McCarthy, Phys. Chem. Chem. Phys., 2019, 21, 2946–2956.
- 64 M. C. McCarthy, K. L. K. Lee, R. A. Loomis, A. M. Burkhardt, C. N. Shingledecker, S. B. Charnley, M. A. Cordiner, E. Herbst, S. Kalenskii, E. R. Willis, C. Xue, A. J. Remijan and B. A. McGuire, Nat. Astron., 2021, 5, 176–180.
- 65 S. J. Goettl, Z. Yang, S. Kollotzek, D. Paul, R. I. Kaiser, A. Somani, A. Portela-Gonzalez, W. Sander, A. A. Nikolayev, V. N. Azyazov and A. M. Mebel, J. Phys. Chem. A, 2023, 127, 5723–5733.
- 66 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, W. Williams, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, Gaussian 16, Revision B.01,

Gaussian, Inc., Wallingford, CT, 2016; see [https://www.](https://www.gaussian.com) [gaussian.com](https://www.gaussian.com).

- 67 A. H. H. Chang, A. M. Mebel, X.-M. Yang, S. H. Lin and Y. T. Lee, J. Chem. Phys., 1998, 109, 2748–2761.
- 68 H. Eyring, S. H. Lin and S. M. Lin, Basic Chemical Kinetics, John Wiley & Sons, Inc., New York, 1980.
- 69 R. I. Kaiser, C. Ochsenfeld, D. Stranges, M. Head-Gordon and Y. T. Lee, Faraday Discuss., 1998, 109, 183–204.
- 70 J. Laskin and C. Lifshitz, J. Mass Spectrom., 2001, 36, 459–478.
- 71 W. B. Miller, S. A. Safron and D. R. Herschbach, Discuss. Faraday Soc., 1967, 44, 108–122.
- 72 R. A. Marcus, J. Chem. Phys., 1952, 20, 359–364.
- 73 D. S. N. Parker, R. I. Kaiser, O. Kostko and M. Ahmed, Chem. Phys. Chem., 2015, 16, 2091–2093.
- 74 C. He, R. I. Kaiser, W. Lu, M. Ahmed, V. S. Krasnoukhov, P. S. Pivovarov, M. V. Zagidullin, V. N. Azyazov, A. N. Morozov and A. M. Mebel, Chem. Sci., 2023, 14, 5369–5378.
- 75 T. Yang, R. I. Kaiser, T. P. Troy, B. Xu, O. Kostko, M. Ahmed, A. M. Mebel, M. V. Zagidullin and V. N. Azyazov, Angew. Chem., Int. Ed., 2017, 129, 4586–4590.
- 76 L. Zhao, R. Kaiser, W. Lu, B. Xu, M. Ahmed, A. N. Morozov, A. M. Mebel, A. H. Howlader and S. F. Wnuk, Nat. Commun., 2019, 10, 1–7.
- 77 L. Zhao, R. I. Kaiser, B. Xu, U. Ablikim, M. Ahmed, M. V. Zagidullin, V. N. Azyazov, A. H. Howlader, S. F. Wnuk and A. M. Mebel, J. Phys. Chem. Lett., 2018, 9, 2620–2626.
- 78 S. J. Goettl, L. B. Tuli, A. M. Turner, Y. Reyes, A. H. Howlader, S. F. Wnuk, P. Hemberger, A. M. Mebel and R. I. Kaiser, J. Am. Chem. Soc., 2023, 145, 15443–15455.