

## CHEMISTRY

# Formation of hydrogen trioxide (HOOOH) in extraterrestrial ice analogs and its role as an oxidizer in prebiotic chemistry

Xilin Bai<sup>1†</sup>, Chunxiao Li<sup>1†</sup>, Yuheng Luo<sup>2†</sup>, Qi'ang Gong<sup>1</sup>, Jinghui Lu<sup>3</sup>, Jiuzhong Yang<sup>3</sup>, Yang Pan<sup>3</sup>, Zhenrong Sun<sup>1</sup>, André K. Eckhardt<sup>4\*</sup>, Rui Sun<sup>2\*</sup>, Ralf I. Kaiser<sup>2\*</sup>, Tao Yang<sup>1\*</sup>

The formation and characterization of hydrogen trioxide have fascinated scientists for more than a century, due to its role as a prototype model for oxygen-chain bonding and as a key transient in antibody-catalyzed oxidation reactions relevant to the origin of life. However, the abiotic formation pathways to hydrogen trioxide have remained elusive. Here, we demonstrate in laboratory simulation experiments that hydrogen trioxide effectively forms in water–molecular oxygen ice analogs at temperatures as low as 5 kelvin under exposure to proxies of galactic cosmic rays. Exploiting synchrotron vacuum ultraviolet photoionization reflectron time-of-flight mass spectrometry, hydrogen trioxide along with hydrogen peroxide and the hydroperoxyl radical was identified during the temperature-programmed desorption of the irradiated ices. This abiotic synthesis expands the oxidant inventory on interstellar grains, icy moons, and Kuiper belt objects, offering a plausible source of essential oxidizers for prebiotic chemistry in space.

## INTRODUCTION

Hydrogen trioxide (HOOOH)—the simplest polyoxide—has long been regarded as a key intermediate in atmospheric (1, 2), biological (3, 4), and environmental processes (1, 5) since its initial proposal by Berthelot (6) in 1880. Extensive investigation of its preparation, reactivity, structural characterization, and quantum chemical properties began only in the 1960s (7). Czapski and Bielski (8) first claimed the production of HOOOH in the pulse radiolysis of acidified air-saturated aqueous solutions, followed by the ultraviolet (UV) absorption spectroscopy of HOOOH in the pulse radiolysis of air-saturated perchloric acid solution (9, 10). Giguère *et al.* (11, 12) identified the fundamental skeletal vibrations of HOOOH via infrared and Raman spectra exploiting electrically discharged mixtures of water (H<sub>2</sub>O), hydrogen peroxide (HOOH), and oxygen (O<sub>2</sub>) at cryogenic temperatures. Plesničar *et al.* (13–16) prepared HOOOH at low temperatures in organic solvents via ozonation of 1,2-diphenylhydrazines or through decomposition of hydrotrioxides (ROOOH) with R being an organic group.

Enlightened by the microscopic reversibility that HOOOH can be reductively generated from ozone (O<sub>3</sub>) decomposition catalyzed by water (17), Wentworth *et al.* (3, 18) suggested that antibodies adopt water as an electron source to catalyze its addition to singlet molecular oxygen (<sup>1</sup>O<sub>2</sub>), yielding HOOOH as a reactive intermediate. Engdahl and Nelander (19) characterized the fundamental vibrations of HOOOH in argon matrices, identifying its antisymmetric O—O stretch at 776 cm<sup>−1</sup> as an ideal marker distinct from strong water absorption. Xu and Goddard *et al.* (1, 20) theoretically delineated plausible mechanisms for HOOOH in a water-catalyzed reaction of water with <sup>1</sup>O<sub>2</sub> or via decomposition of [(HO<sub>2</sub>)(HO<sub>3</sub>)] hydrogen-bonded complexes. Suma *et al.* (21) reported the first pure rotational

spectra and precise geometrical structure of *trans*-HOOOH. High-level ab initio studies including the benchmark CCSD(T)-F12 calculations by Hollman and Schaefer (22) confirmed a zigzag skew-chain helical-type structure of *trans*-HOOOH with C<sub>2</sub> symmetry.

Despite the astrochemical and astrobiological importance of HOOOH in the interstellar medium and the Solar System, its formation mechanism under extraterrestrial-like conditions has remained poorly understood. Water ice—a potential precursor to HOOOH—is found ubiquitously in cold molecular clouds (23, 24) and throughout the Solar System including Kuiper belt objects (25, 26), comets (27, 28), and icy moons such as Europa, Ganymede, and Callisto (29, 30). Molecular oxygen—a critical molecular building block of HOOOH—has been detected in molecular clouds, e.g., Orion (31) and ρ Ophiuchi (Oph) A (32), and the condensed phase on the Galilean and Saturnian moons (33, 34). Notably, molecular oxygen also exhibits high abundances in the coma of comet 67P/Churyumov-Gerasimenko (35).

The coexistence of H<sub>2</sub>O and O<sub>2</sub> in interstellar and Solar System ices suggests that HOOOH could form via nonequilibrium processes such as through the interaction of H<sub>2</sub>O–O<sub>2</sub> ices with ionizing radiation such as galactic cosmic rays (GCRs) or electrons trapped in the magnetospheres of the giant planets. However, experimental studies on irradiated H<sub>2</sub>O–O<sub>2</sub> ices are limited. Under conditions simulating the icy Galilean moons, Cooper, Moore, and Hudson detected HOOH in proton irradiated H<sub>2</sub>O and H<sub>2</sub>O–O<sub>2</sub> ices (36) along with the hydroperoxyl (HOO) and hydrotrioxyl (HOOO) radicals (37, 38). Zheng *et al.* (39) conducted electron irradiation studies of H<sub>2</sub><sup>18</sup>O–O<sub>2</sub> ices and tentatively assigned the formation of HOOOH via a single HOO bending mode at 1350 cm<sup>−1</sup>. However, the symmetric/antisymmetric HOO bending modes of HOOOH at 1347 and 1359 cm<sup>−1</sup> overlap strongly with water absorption bands, thus complicating unambiguous identification (19). Consequently, conclusive evidence for the formation and identification of HOOOH in simulated extraterrestrial environments still remains elusive.

In this study, we present persuasive evidence on the formation and identification of HOOOH, alongside HOO, HOOH, and O<sub>3</sub> in low-temperature H<sub>2</sub>O–O<sub>2</sub> ice analogs exposed to proxies of GCRs.

<sup>1</sup>State Key Laboratory of Precision Spectroscopy, East China Normal University, Shanghai 200062, P. R. China. <sup>2</sup>Department of Chemistry, University of Hawaii at Manoa, Honolulu, HI 96822, USA. <sup>3</sup>National Synchrotron Radiation Laboratory, University of Science and Technology of China, Hefei, Anhui 230029, P. R. China. <sup>4</sup>Lehrstuhl für Organische Chemie II, Ruhr-Universität Bochum, Bochum 44801, Germany.

\*Corresponding author. Email: tyang@ps.ecnu.edu.cn (T.Y.); ralfk@hawaii.edu (R.I.K.); ruisun@hawaii.edu (R.S.); andre.eckhardt@ruhr-uni-bochum.de (A.K.E.)

†These authors contributed equally to this work.

By exploiting the temperature-programmed desorption (TPD) technique, we investigate the subliming products formed in the irradiated ices via the synchrotron vacuum UV photoionization reflectron time-of-flight mass spectrometry (SVUV-PI-ReTOF-MS) (40–42). We achieve unambiguous identification of HOOOH plus other key oxidants through photoionization efficiency (PIE) curve fitting, in conjunction with high-level quantum chemical calculations, quadrupole mass spectrometry, and isotopic labeling experiments. Our results demonstrate that HOOOH can form effectively on icy grains and surfaces of Galilean and Saturnian moons and comet 67P/Churyumov-Gerasimenko. These findings expand the molecular inventory of oxidants in space and enhance our understanding of their roles as key chemical sources to oxidize complex organics or even fuel extraterrestrial life in the outer Solar System. Our work may also motivate astronomical searches for HOOOH in star-forming regions and on the icy moons of Jupiter and Saturn, where HOOH has already been identified.

## RESULTS

### Computations

The SVUV-PI-ReTOF-MS method represents a unique technique to identify subliming species in the TPD phase based on the mass-to-charge ratio ( $m/z$ ) of the ions and the adiabatic ionization energies (IEs) of the neutrals (40, 41). Because experimental IEs of  $\text{H}_2\text{O}_3$  isomers are unknown, we theoretically determined their adiabatic IEs at a level of sufficient accuracy (0.01 to 0.08 eV) (43). Our computations revealed two structural isomers, with geometries calculated at the CCSD(T)/aug-cc-pVTZ level of theory and adiabatic IEs and relative energies computed at the CCSD(T)/CBS//B3LYP/aug-cc-pVTZ level of theory (Fig. 1 and tables S1 to S3). Two conformers of HOOOH were located: *trans*-HOOOH ( $C_2$  symmetry) and *cis*-HOOOH ( $C_s$  symmetry); these conformers differ only by 9 kJ mol<sup>−1</sup> and are interconvertible via a low *cis-trans* isomerization barrier of only 22 kJ mol<sup>−1</sup> (5, 44, 45). The hyperconjugative interaction between the lone pair of electrons at the oxygen atoms and the  $\sigma^*$  orbital of the adjacent O–H bond optimizes the dihedral angle, thus achieving maximum stability and electron delocalization. Specifically, the dihedral angle for *trans*-HOOOH is 81.3°, whereas the corresponding angle for HOOH is 112° (5). This hyperconjugation effect also shortens the O–O bond length in *trans*-HOOOH to 143.5 pm, in contrast to the typical O–O bond length of 145.6 pm in HOOH. We also identified the water-oxygen molecular complex  $[\text{O}(\text{H}_2\text{O})\text{O}]$  isomer that is less stable by 485 kJ mol<sup>−1</sup> compared to *trans*-HOOOH. This  $C_{2v}$  symmetric isomer holds a central oxygen atom with a  $sp^3$  hybridization,

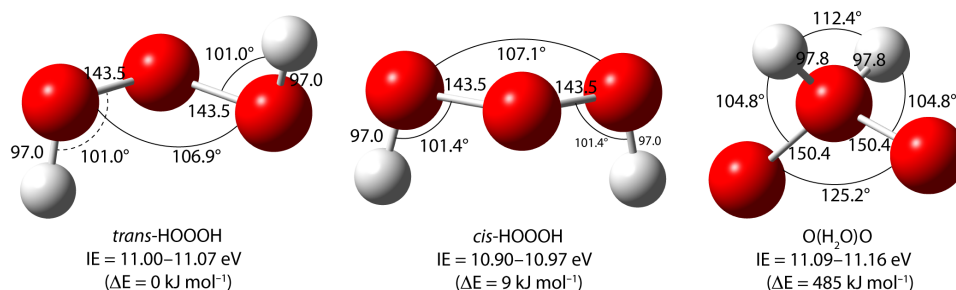
resulting in one unpaired electron on each of the terminal oxygen atoms. The adiabatic IEs of the aforementioned isomers were computed respectively to be 11.00 to 11.07 eV (*trans*-HOOOH), 10.90 to 10.97 eV (*cis*-HOOOH), and 11.09 to 11.16 eV  $[\text{O}(\text{H}_2\text{O})\text{O}]$  (5, 45).

### Mass spectrometry

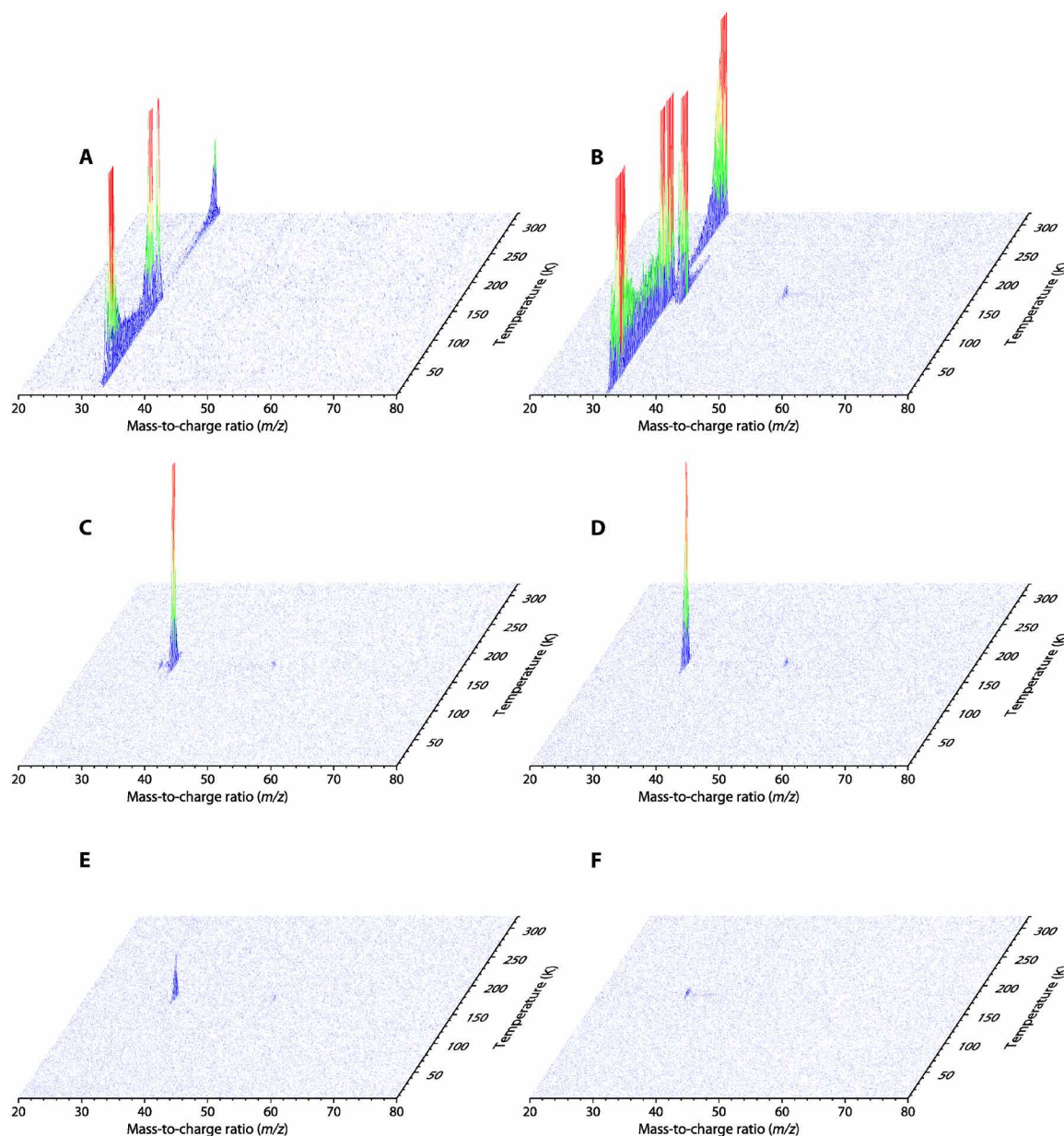
In the TPD phase of irradiated  $\text{H}_2\text{O}-\text{O}_2$  ices, ion signals were observed at  $m/z$  of 32, 33, 34, and 50 at photon energies of 12.20 eV (Figs. 2 and 3, A to B). Because the products only contain hydrogen and oxygen atoms, we can link them to the molecular formulas of  $\text{O}_2$  (32 u),  $\text{HO}_2$  (33 u),  $\text{H}_2\text{O}_2$  (34 u), and  $\text{H}_2\text{O}_3$  (50 u), respectively. The blank experiment, which was conducted without exposing the ices to ionizing radiation, reveals only signal at  $m/z$  of 32 ( $\text{O}_2^+$ ), i.e., the ionized molecular oxygen reactant, suggesting that  $\text{HO}_2$ ,  $\text{H}_2\text{O}_2$ , and  $\text{H}_2\text{O}_3$  are the result of the radiation exposure of the  $\text{H}_2\text{O}-\text{O}_2$  ices.

Isotopic labeling experiments involving the  $\text{H}_2^{18}\text{O}-\text{O}_2$  ices provided additional evidence on the formation of isotopically substituted products (Fig. 3, C and D, and figs. S1 and S2). For clarity, we distinguish the oxygen atoms in the reactants and products as  $^{16}\text{O}$  and  $^{18}\text{O}$ , respectively. Here, TPD profiles at  $m/z$  of 50, 52, 54, and 56 in the  $\text{H}_2^{18}\text{O}-^{16}\text{O}_2$  ice at the photon energy of 12.20 eV can be respectively assigned to ion signals of  $\text{H}_2^{16}\text{O}_3$ ,  $\text{H}_2^{18}\text{O}^{16}\text{O}_2$ ,  $\text{H}_2^{18}\text{O}_2^{16}\text{O}$ , and  $\text{H}_2^{18}\text{O}_3$ , which yield a branching ratio of 1:1.25:1.03:1.13 after the integration of each TPD profile at distinct  $m/z$ . Similarly, TPD profiles at  $m/z$  of 33, 35, and 37 at the photon energy of 11.40 eV correspond to ions of  $\text{H}^{16}\text{O}_2$ ,  $\text{H}^{16}\text{O}^{18}\text{O}$ , and  $\text{H}^{18}\text{O}_2$  whose branching ratio is 1:0.74:1.07, while those at  $m/z$  of 34, 36, and 38 at the photon energy of 11.40 eV correspond to ions of  $\text{H}_2^{16}\text{O}_2$ ,  $\text{H}_2^{16}\text{O}^{18}\text{O}$ , and  $\text{H}_2^{18}\text{O}_2$  whose branching ratio is 1:1.24:0.90.

We also analyzed the subliming molecules via an electron-impact quadrupole mass spectrometer (EI-QMS) operating with 70-eV electrons. Here, ion signals were detected at  $m/z$  of 34 and 48 in the  $\text{H}_2^{16}\text{O}-^{16}\text{O}_2$  ices, indicating the formation of  $\text{H}_2^{16}\text{O}_2$  and  $^{16}\text{O}_3$  (fig. S3). We cannot assign ion signals at  $m/z$  of 33 to  $\text{H}^{16}\text{O}_2^+$  due to strong mass interference of  $^{16}\text{O}_2^+$ . The absence of signals at  $m/z$  of 50 ( $\text{H}_2^{16}\text{O}_3^+$ ) can be attributed to low ion counts and limited sensitivity of EI-QMS. In the irradiated  $\text{H}_2^{18}\text{O}-^{16}\text{O}_2$  ices, we may assign ion signals at  $m/z$  of 34, 36, and 38 as  $\text{H}_2^{16}\text{O}_2^+/\text{H}^{16}\text{O}^{18}\text{O}^+$ ,  $\text{H}_2^{16}\text{O}^{18}\text{O}^+/\text{H}^{18}\text{O}_2^+$ , and  $\text{H}_2^{18}\text{O}_2^+$ , respectively (fig. S4). Ion signals at  $m/z$  of 33, 35, and 37 can originate from either mass interference or fragmentation of neighboring ions. We detected ions at  $m/z$  of 48, implying preferential formation of  $^{16}\text{O}_3$  over other  $\text{O}_3$  isotopologues (39). Similarly, we did not observe any  $\text{H}_2\text{O}_3$  isotopologues via EI-QMS in the irradiated  $\text{H}_2^{18}\text{O}-^{16}\text{O}_2$  ices.



**Fig. 1. CCSD(T)/aug-cc-pVTZ optimized geometries of  $\text{H}_2\text{O}_3$  isomers.** Bond lengths and angles are in the units of picometers (pm) and degrees (°), respectively, while relative energies with respect to *trans*-HOOOH ( $\Delta E$ ) are given in the units of kilojoules per mole (kJ mol<sup>−1</sup>). Corresponding adiabatic IEs are provided in electron volts (eV).



**Fig. 2. SVUV-PI-ReTOF-MS at selected photon energies.** SVUV-PI-ReTOF-MS data recorded during the TPD phase of nonirradiated H<sub>2</sub>O-O<sub>2</sub> ice at the photon energy (PE) of (A) 12.20 eV (blank experiment) and irradiated H<sub>2</sub>O-O<sub>2</sub> ice at PEs of (B) 12.20 eV, (C) 11.40 eV, (D) 11.30 eV, (E) 11.00 eV, and (F) 10.90 eV, respectively.

### PIE curves

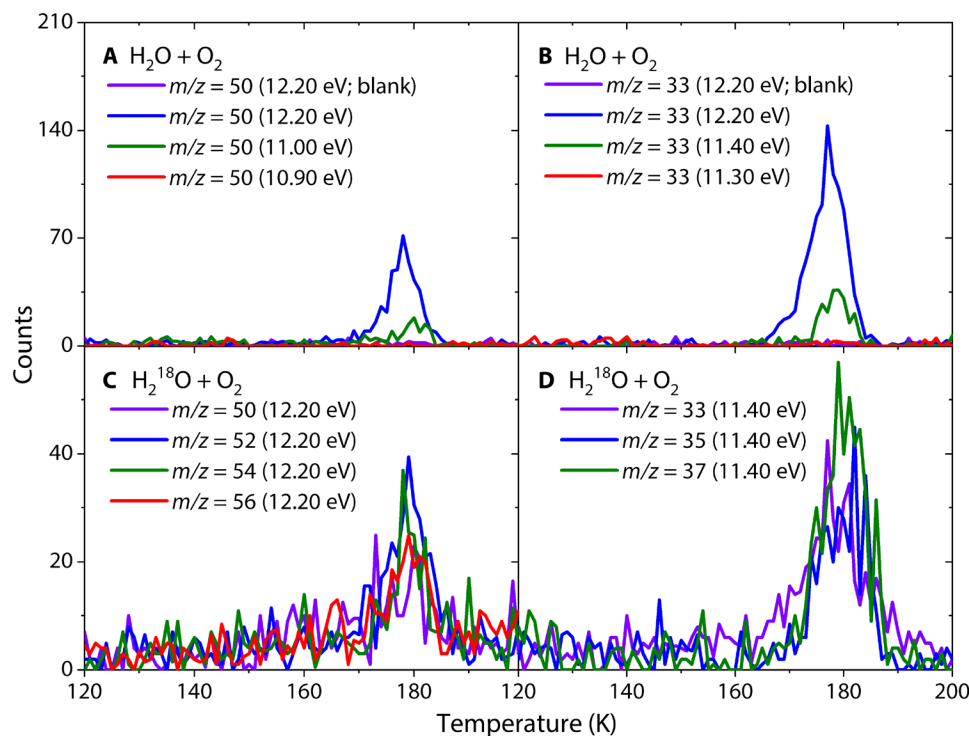
The PIE curves report the intensity of ion counts versus the photon energy at distinct  $m/z$ . For the irradiated H<sub>2</sub>O-O<sub>2</sub> ices, the experimental PIE curves of photoionized products at  $m/z$  of 33 (HO<sub>2</sub><sup>+</sup>), 34 (H<sub>2</sub>O<sub>2</sub><sup>+</sup>), and 50 (H<sub>2</sub>O<sub>3</sub><sup>+</sup>) were collected by scanning the photon energies from 10.80 to 12.20 eV in an interval of 0.10 eV (HO<sub>2</sub><sup>+</sup> and H<sub>2</sub>O<sub>3</sub><sup>+</sup>) and from 10.40 to 11.60 eV in an interval of 0.05 eV (H<sub>2</sub>O<sub>2</sub><sup>+</sup>), respectively, in the temperature range of 175 to 184 K (Fig. 4, A to C, and fig. S5). These PIE curves at  $m/z$  of 33, 34, and 50 display the onsets of ion counts at  $11.35 \pm 0.05$  eV,  $10.60 \pm 0.03$  eV, and  $10.95 \pm 0.05$  eV, respectively.

The reported IE of  $11.352 \pm 0.007$  eV for the hydroperoxyl radical HOO (46) matches well with the experimental onset of ion

counts for  $m/z$  of 33 at  $11.35 \pm 0.05$  eV, whereas the IE of HOOH of  $10.631 \pm 0.007$  eV (46) correlates nicely with the onset of ion counts for  $m/z$  of 34 at  $10.60 \pm 0.03$  eV. We emphasize that the appearance energy of HO<sub>2</sub><sup>+</sup> fragment from HOOH was determined to be  $15.112 \pm 0.035$  eV (46), which cannot contribute to our experimental PIE curve recorded at  $m/z$  of 33. In addition, PIE curves collected in the present experiments at  $m/z$  of 33 and 34 agree exceptionally well with known literature PIE curves of HOO (46) and HOOH (47), respectively.

The onset of ion counts for  $m/z$  of 50 (H<sub>2</sub>O<sub>3</sub><sup>+</sup>) at  $10.95 \pm 0.05$  eV is in line with the calculated IEs of *trans*-HOOOH (11.00 to 11.07 eV) and *cis*-HOOOH (10.90 to 10.97 eV), indicating the formation of *trans*-/*cis*-HOOOH, but not O(H<sub>2</sub>O)O (11.09 to 11.16 eV). Nevertheless, no





**Fig. 3. TPD profiles.** TPD profiles for (A)  $m/z = 50$  of nonirradiated  $\text{H}_2\text{O}-\text{O}_2$  ice at the PE of 12.20 eV (blank experiment) and irradiated  $\text{H}_2\text{O}-\text{O}_2$  ices at PEs of 12.20, 11.00, and 10.90 eV, and (B)  $m/z = 33$  of nonirradiated  $\text{H}_2\text{O}-\text{O}_2$  ice at the PE of 12.20 eV and irradiated  $\text{H}_2\text{O}-\text{O}_2$  ices at PEs of 12.20, 11.40, and 11.30 eV, respectively. TPD profiles obtained at (C)  $m/z = 50, 52, 54$ , and  $56$  at the PE of 12.20 eV and (D)  $m/z = 33, 35$ , and  $37$  at the PE of 11.40 eV of irradiated  $\text{H}_2^{18}\text{O}-\text{O}_2$  ices.

reference PIE curve of any  $\text{H}_2\text{O}_3$  isomer is available. Our computations in Franck-Condon factors (FCFs) from *trans*-/*cis*- $\text{HOOOH}^+$  to the resulting cations including *trans-trans*- $\text{HOOOH}^+$ , *cis-trans*- $\text{HOOOH}^+$ , and *cis-cis*- $\text{HOOOH}^+$  (fig. S6) revealed that FCFs for the ionization from *trans*- $\text{HOOOH}$  to *cis-cis*- $\text{HOOOH}^+$  are at least two orders of magnitude higher than the others. The computations correlate with the rapidly increasing PIE profile at an  $m/z$  of 50, which indicates an efficient adiabatic ionization at  $10.95 \pm 0.05$  eV (Fig. 4C). This PIE curve then drops sharply in intensity beyond  $11.25 \pm 0.05$  eV, suggesting an efficient dissociative ionization.

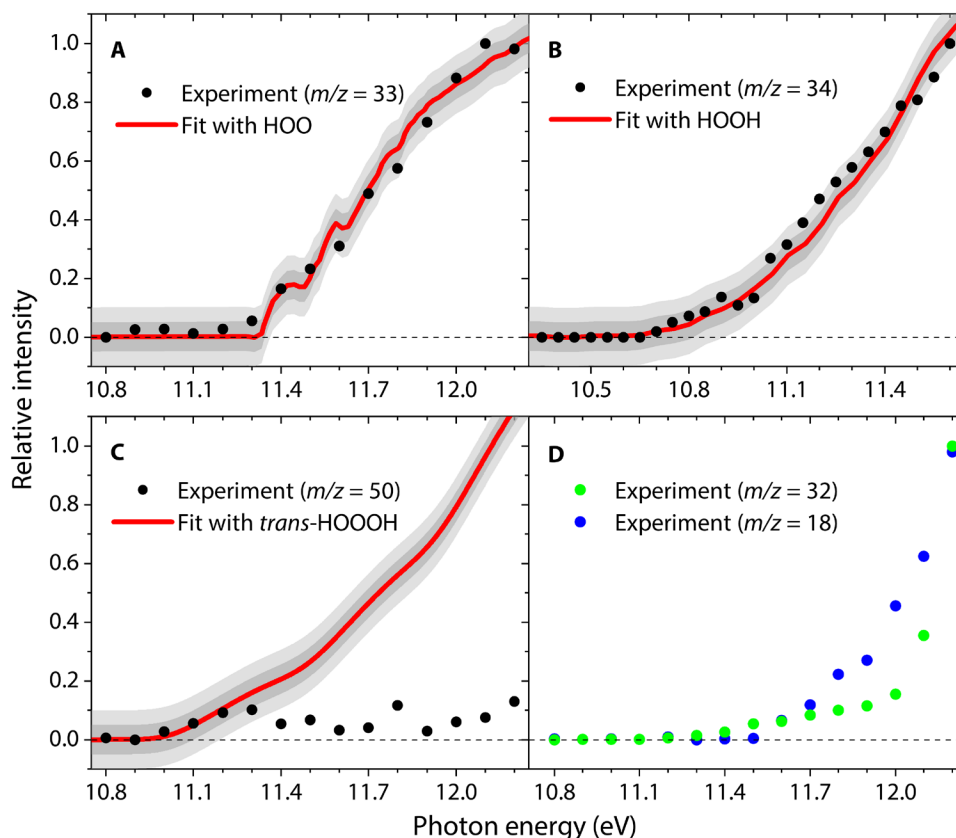
## DISCUSSION

We further conducted an extensive computational search of entrance and exit channels leading to  $\text{HOOOH}$  on the global potential energy surface (PES) involving  $\text{HOOOH}$  neutrals and cations to elucidate the confirmative detection of  $\text{HOOOH}$  (Fig. 5). The computations predict that *cis-cis*- $\text{HOOOH}^+$  isomerizes promptly to the *cis-trans*- $\text{HOOOH}^+$ , which undergoes unimolecular decomposition via a barrier of 1.09 eV to  $\text{H}_2\text{O}$  plus  $\text{O}_2^+$  at the CCSD(T)/CBS//B3LYP/aug-cc-pVTZ + zero-point vibrational energy (ZPVE) level of theory. To account for potential multireference characters of this transition state, the dissociation barrier was further explored at the complete active-space second-order perturbation theory CASPT2(15,13)/def2-TZVP//B3LYP/aug-cc-pVTZ + ZPVE level of theory. At this level, the barrier turns out to be 11.30 eV and matches well the experimental PIE drop at  $11.25 \pm 0.05$  eV for  $m/z$  of 50. In addition, this barrier also coincides with the appearance energy of  $11.25 \pm 0.05$  eV in the dissociative channel of  $\text{H}_2\text{O}$  plus  $\text{O}_2^+$  as indicated in the

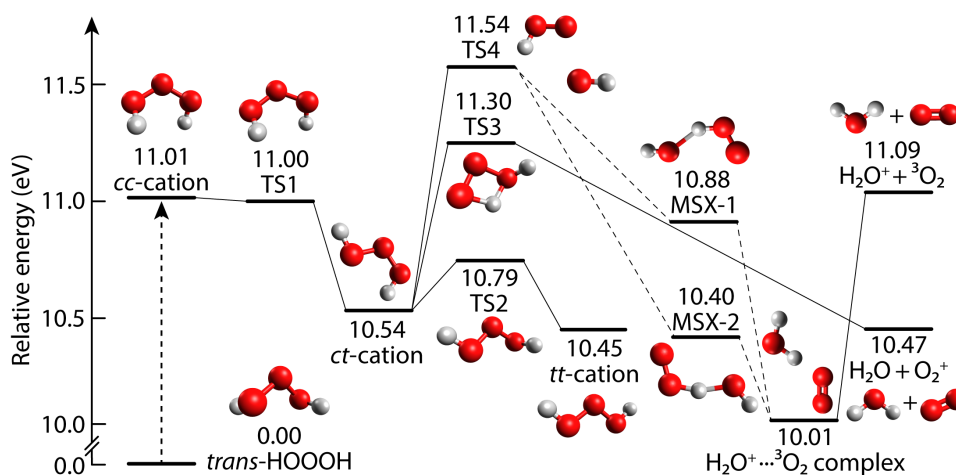
experimental PIE curve of  $\text{O}_2^+$  (Fig. 4D and fig. S7). Notice that the ionization threshold of  $\text{O}_2$  itself was measured to be  $12.071 \pm 0.001$  eV (48).

We also carried out the comprehensive analysis of potential intersystem crossings (ISCs) in the reaction pathway involving the dissociative channel  $\text{H}_2\text{O}^+$  plus  $^3\text{O}_2$  (fig. S8). Starting from *cis-trans*- $\text{HOOOH}^+$ , the CASPT2 energy scans showed that the energy on the doublet manifold increases, while the quartet state energy decreases as the  $\text{HOO}-\text{OH}$  bond stretches, resulting in the identification of the transition state TS4 at 11.54 eV. We located two minimum-energy crossing points MSX-1 and MSX-2 connecting to this transition state between the doublet and quartet surfaces, either of which can lead to the  $\text{H}_2\text{O}^+ \cdots ^3\text{O}_2$  complex that dissociates barrierlessly to  $\text{H}_2\text{O}^+$  plus  $^3\text{O}_2$ . This dissociative pathway could be responsible for the ionization onset of  $11.55 \pm 0.05$  eV in the experimental PIE curve of  $\text{H}_2\text{O}^+$  (Fig. 4D and fig. S7); this value is lower than the ionization threshold of  $\text{H}_2\text{O}$  located at  $12.619 \pm 0.006$  eV (49). Alternative dissociative pathways of  $\text{HOOOH}^+$  leading to  $\text{H}_2\text{O}^+$  plus  $^1\text{O}_2$ ,  $\text{HO}_2^+$  plus OH,  $\text{HO}_3^+$  plus H, and  $\text{HO}_2$  plus  $\text{OH}^+$  are located 12.32, 12.98, 14.44, and 16.87 eV higher than *trans*- $\text{HOOOH}$ , respectively, and hence are closed energetically in our photoionization studies (tables S1 to S4). Overall, the agreement on the PIE curves and the ionization onsets of  $\text{HOOOH}$  along with the dissociative channels of  $\text{H}_2\text{O}$  plus  $\text{O}_2^+$  and  $\text{H}_2\text{O}^+$  plus  $^3\text{O}_2$  between experiments and computations strongly supports the formation and detection of  $\text{HOOOH}$  molecules in our irradiated  $\text{H}_2\text{O}-\text{O}_2$  ices.

Is it possible to detect  $\text{HOOOH}$  in deep space? Notably,  $\text{HOO}$  and  $\text{HOOH}$ —two coproducts of the interaction of ionizing radiation with  $\text{H}_2\text{O}-\text{O}_2$  ices in the present experiments—have been detected



**Fig. 4. PIE curves.** PIE curves for the species linked to (A)  $m/z = 33$ , (B)  $m/z = 34$ , (C)  $m/z = 50$ , and (D)  $m/z = 18$  and  $32$ . The black, blue, and green dots represent the experimental data, while the red lines indicate the calculated or reference PIE curves. The experimental PIE curve at  $11.25 \pm 0.05$  eV, while the experimental PIE curves at  $m/z = 32$  and  $18$  rise at  $11.25 \pm 0.05$  and  $11.55 \pm 0.05$  eV, respectively. The uncertainties in the fit are noted with dark-gray ( $1\sigma$ ) and light-gray ( $2\sigma$ ) shaded regions.



**Fig. 5. The PES of the HOOOH neutral, cations, and dissociative channels.** The *tt*-/ *ct*-/ *cc*-cation stands for the radical cations *trans-trans*-/ *cis-trans*-/ *cis-cis*-HOOOH<sup>+</sup>, respectively. The dashed arrow indicates the adiabatic ionization process from *trans*-HOOOH to *cc*-cation, while the dashed lines connecting two minimum-energy crossing points (MSX-1 and MSX-2) represent the pathways involving the intersystem crossing (ISC). Atoms are color coded in white (hydrogen) and red (oxygen).

toward the cloud complex  $\rho$  Oph A (50, 51).  $\rho$  Oph A also represents one of two interstellar clouds in which gas-phase  $O_2$  was detected with abundance ratios of  $HOO/O_2 \approx HOOH/O_2 \approx 0.6 \times 10^{-3}$  (32, 52). Since both  $H_2O$  and  $O_2$  are present toward  $\rho$  Oph A and the coproducts  $HOO$  and  $HOOH$  have been also detected,  $HOOOH$  should be existing after formation and sublimation from icy grains. On the other side, the abundance ratios of  $HOO$  and  $HOOH$  versus  $O_2$  in the coma of comet 67P/Churyumov-Gerasimenko as determined by the Double Focusing Mass Spectrometer of the ROSINA instrument, are very close to those interstellar ones (53). The in situ measurement revealed local abundances of  $O_2$  relative to  $H_2O$  ranging from 1 to 10%, suggesting that  $O_2$  is chemically relevant to  $H_2O$  ice by radiolysis rather than  $CO_2$  or  $CO$  ices (35). Therefore, astrophysical  $H_2O$ - $O_2$  ices under exposure of GCRs can be regarded as chemical and energetic sources in producing oxidizing agents including  $HOOOH$ ,  $HOOH$ ,  $HOO$ , and  $O_3$  relevant to the extraterrestrial origin of life in deep space and the outer Solar System. The rotational frequency of *trans*- $HOOOH$  (20838.6 MHz) (21) is potentially linked to the unidentified line (20838.2 MHz) observed in the W51 region (54), a massive molecular complex containing clusters of  $H_2O$  masers.

In conclusion, the present work provides compelling evidence for the identification of key oxidizing agents  $HOOOH$ ,  $HOOH$ ,  $HOO$ , and  $O_3$  prepared in water-molecular oxygen ( $H_2O$ - $O_2$ ) ice analogs upon exposure to ionizing radiation in the form of proxies of GCRs. Although  $HOOOH$  has been studied extensively for more than one century, we demonstrate the formation and detection of  $HOOOH$  in interstellar and Solar System analog ices exploiting state-of-the-art SVUV-PI-ReTOF-MS, in association with high-level quantum chemical calculations, quadrupole mass spectrometry, and isotopic labeling experiments. The laboratory experiments conducted under conditions simulating astrophysical ones suggest that  $HOOOH$  can be produced in ices rich in  $H_2O$  and  $O_2$  under the influence of ionizing radiation. The studies further suggest that  $HOOOH$  should be searched for and detected in the gas phase after sublimation from the icy grains toward  $\rho$  Oph A. The detection of interstellar  $HOOOH$  combined with updated astrochemical modeling would provide fundamental constraints and knowledge on the chemistry of oxygen-containing inorganic molecules in deep space. It would further provide a hitherto overlooked entry point to key oxidizers in prebiotic chemistry, eventually deciphering prebiotic redox cycles that could fuel microbial life on icy moons (55) or even oxygen-driven catalytic cycles such as the antibody catalysis in the oxidation of water (56).

## MATERIALS AND METHODS

### Experimental

The surface scattering experiments were conducted at the Shanghai-Hawaii-Hefei Advanced Research Center, using the VUV synchrotron radiation at the Combustion and Flame Beamline of the National Synchrotron Radiation Laboratory (NSRL) (40, 41). The surface machine consists of a hydrogen-free stainless steel ultrahigh vacuum chamber that can be pumped down to a few  $10^{-11}$  torr. The ice analogs were deposited onto a mirror-polished silver substrate ( $12.6 \times 15.1$  mm<sup>2</sup>) attached to an oxygen-free high thermal conductivity copper coldfinger at  $4.8 \pm 0.1$  K. High-purity  $H_2O$  (Macklin, high-performance liquid chromatography grade) and molecular  $O_2$  (Air Liquide; 99.999%) were introduced and condensed onto the silver substrate via two sets of leak valves and glass capillary arrays (10-mm array diameter) at pressures of  $(1 \pm 0.1) \times 10^{-8}$  and

$(2 \pm 0.2) \times 10^{-8}$  torr, respectively, achieving an actual depositing ratio ( $[H_2O]/[O_2]$ ) of  $6.5 \pm 1.0:1$ . This ratio was determined by an EI-QMS (Pfeiffer, QMG 220 M2). Before experiments using  $H_2^{18}O$  (Aladdin; 97.0 atom %  $^{18}O$ ), tubings and chambers were purged with  $H_2^{18}O$  vapor to allow any  $^{16}O/^{18}O$  isotopic exchange to occur before the preparation of  $H_2^{18}O$ - $O_2$  ices. The thickness of deposited ices was measured via laser interferometry with a HeNe 632.8-nm laser (Thorlabs, HNL008L), which was determined to be  $755 \pm 50$  nm based on analyzing the interference patterns and the average refractive index of the two components in amorphous ices,  $1.26 \pm 0.02$  (Supplementary Methods and fig. S9) (57, 58).

After deposition, the silver substrate was rotated and translated to the irradiation position and subjected to isothermal irradiation with 5-keV electrons (SPECS GmbH, EQ 22/35 electron source) at an incident angle of  $70^\circ$  relative to the surface normal. Distinct experiments were carried out under blank (0 nA for 60 min) and irradiated (1000 nA for 60 min, 1000 nA for 10 min, and 100 nA for 60 min) conditions (fig. S10). Using the CASINO 2.42 simulation software (59), the average penetration depth of the electrons (1000-nA current) was calculated to be  $226 \pm 23$  nm (fig. S11), significantly less than the ice thickness, ensuring no interaction between the electrons and the silver substrate. The highest irradiation doses were calculated to be 99.2 eV per  $H_2O$  and 176.4 eV per  $O_2$ , respectively (table S5).

After irradiation, the ice sample was repositioned for TPD, annealing from 5 to 320 K at a rate of  $1 \text{ K min}^{-1}$  controlled and monitored by a programmable temperature controller (Lake Shore, Model 336) in an accuracy of  $\pm 0.1$  K. Simultaneously, subliming molecules were mass analyzed via SVUV-PI-ReTOF-MS. Molecular ions at distinct  $m/z$  generated upon the selected photon energy were extracted by pulsed grid voltage ( $-150$  V for 2.5  $\mu$ s) and detected by a 40-mm dual-multichannel plate detector (Jordan TOF Products, C-726), whose signals were preamplified (AMETEK Inc., ORTEC 9306), discriminated and amplified (Advanced Research Instruments Corp., F-100TD), and recorded by a multichannel scaler (FAST ComTec, P7889) with 3.2-ns accuracy. Typical ReTOF mass spectra were generated at an experimental repetition rate of 15 kHz and a data acquisition time of 60 s, in accordance with the TPD rate. In the TPD range between 175 and 184 K, the SVUV photon energy was tuned from 10.80 to 12.20 eV in an interval of 0.10 eV or from 10.40 to 11.60 eV in an interval of 0.05 eV, to extract PIE curves of targeted molecules, which report the intensities of ions at specific  $m/z$  as a function of photon energy. During the SVUV photon energy scan, the obtained TPD spectra were normalized with respect to the TPD spectrum at 12.20 or 11.60 eV and the corresponding photon flux. Subliming molecules were also detected during TPD via the EI-QMS in the residual gas analyzer mode.

### Computational

All density functional theory computations were carried out with Gaussian 16, Revision C.01 (60). For geometry optimizations and frequency computations, the B3LYP functional (61) was used with the Dunning correlation consistent split-valence basis set aug-cc-pVTZ (62). Vibrational frequency calculations were used to characterize stationary points and to obtain ZPVE corrections. Minimum structures were characterized by zero imaginary frequencies, while transition states exactly showed one imaginary frequency. On the basis of the minimum geometries, the corresponding frozen-core coupled cluster (63) CCSD(T)/aug-cc-pVTZ and CCSD(T)/aug-cc-pVQZ

single-point energies were computed and extrapolated to complete basis set limit (64) CCSD(T)/CBS using the built-in extrapolation routine in ORCA 5.0.4 (65). The adiabatic IEs were computed by taking the electronic CCSD(T)/CBS energies and the B3LYP/aug-cc-pVTZ ZPVE-corrected energy difference between the neutral and ionic species that correspond to similar conformations.

The analysis of ISCs between doublet and quartet states was carried out on a set of 49 smoothly transitioning structures along this pathway using Molpro (version 2025.1) (66, 67). Calculations were performed using state-averaged complete active-space self-consistent field (68) wave functions, followed by CASPT2 (69–71) with the def2-TZVP basis set (72). The active space consisted of nine 2p orbitals from oxygen atoms and two 1s orbitals from hydrogen atoms. Two lowest doublet states and one lowest quartet state were considered to avoid root-flipping issues associated with nearly degenerate doublet states. A level shift of 0.3 was applied during CASPT2 calculations to prevent intruder state problems as recommended by Roos and Andersson (73).

## Supplementary Materials

This PDF file includes:

Supplementary Methods

Figs. S1 to S11

Tables S1 to S5

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